

Introduction

The solution for the differential equation is: Topological photonics has emerged as a frontier in controlling light propagation, with topological insulators playing a pivotal role in trapping light at specific frequencies. Traditionally, the design of these materials involves finding the right material composition to achieve the desired optical properties. Our method employs the transfer matrix method to solve the direct problem, which involves understanding how various topological parameters, such as those of We will have different σ_A and σ_B for different materials A and B, θ is the materials, affect the frequencies at which light is fraction of each unit period with material A. trapped. By mapping out these relationships, our study then proceeded to tackle the inverse challenge: determining the combination of two distinct materials necessary to fabricate a topological insulator capable of confining light at a predetermined frequency.

Methods

Wave Propagation

We used the Helmholtz Equation to determine the behavior of a wave in a periodically layered halfspace. Since the medium is invariant in the x direction, we assume the solution to be of the form $u(z)e^{i\xi x}$, where u(z) satisfies

$$u(z)'' + \frac{\omega^2}{c(z)^2}u(z) - \xi^2 u(z) = 0$$

 ω : angular frequency c: speed of light ξ : spatial wave number



Original Position z

To model the periodic medium, we assume. We normalize $z \to z'p, \omega \to \omega'p, \xi \to \xi'p$ (and drop the primes)

Inverse Problem for Topological Photonics

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Propagation Matrix

Let
$$\sigma = \sqrt{\frac{\omega^2}{c^2} - \xi^2}$$

 $u = A\cos(\sigma z) + B\sin(\sigma z)$
 $u' = -\sigma A\sin(\sigma z) + \sigma B\cos(\sigma z)$
 $\begin{bmatrix} u \\ u' \end{bmatrix} (0) = \begin{bmatrix} A \\ \sigma B \end{bmatrix}$
 $\begin{bmatrix} u \\ u' \end{bmatrix} (z) = \begin{bmatrix} \cos(\sigma z) & \frac{\sin(\sigma z)}{\sigma} \\ -\sigma\sin(\sigma z) & \cos(\sigma z) \end{bmatrix} \begin{bmatrix} u \\ u' \end{bmatrix} (0)$

$$T_A = \begin{bmatrix} \cos(\sigma_A \theta) \\ -\sigma_A \sin(\sigma_A \theta) \end{bmatrix}$$

$$T_B = \begin{bmatrix} \cos(\sigma_B(1-\theta)) \\ -\sigma_B \sin(\sigma_B(1-\theta)) \end{bmatrix}$$

Transfer Matrix (In our model, we only consider one layer of A and one layer of B in a period, without refraction) Then the transfer matrix is defined as, with z' = np:

 $T^{(1)} = (T_A T_B)^n$

Results (Direct Problem)

Edge States After calculating the transfer matrix, the edge states occur when the imaginary part of eigenvalues of $T^{(1)}$ is 0 (eigenvalues of $T^{(1)}$ are real).



$$\frac{\sin(\sigma_A \theta)}{\sigma_A}$$
$$\cos(\sigma_A \theta)$$

$$\frac{\sin(\sigma_B(1-\theta))}{\sigma_B}$$
$$\cos(\sigma_B(1-\theta))$$



Visualization of wave propagation with $\xi = 2, \theta = 0.3$

Given a desired frequency ω_D and desired wave number ξ_D , find a structure that supports and edge mode at these parameters. That means finding c_A , c_{R}, θ that meets this design objective.

stated methods.

- Create datasets of edge modes at different parameter values.
- Develop a machine learning algorithm that solves the inverse design problem.

Pilozzi, Laura, et al. "Machine Learning Inverse Problem for Topological Photonics." Communications Physics, vol. 1, no. 57, 2018. https://doi.org/10.1038/s42005-018-0058-8.

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Results (Edge Modes)

Inverse Design Problem

Future Steps

- Compute the direct problem efficiently using

References